

# Zero forcing, linear and quantum controllability for systems evolving on networks

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## Abstract

We study the dynamics of systems on networks from a linear algebraic perspective. The control theoretic concept of *controllability* describes the set of states that can be reached for these systems. Under appropriate conditions, there is a connection between the quantum (Lie theoretic) property of controllability and the (classical) linear systems controllability condition. We investigate how the graph theoretic concept of a zero forcing set impacts the controllability property. In particular, we prove that if a set of vertices is a zero forcing set, the associated dynamical system is controllable. The results open up the possibility of further exploiting the analogy between networks, linear control systems theory, and quantum systems Lie algebraic theory. This study is motivated by several quantum systems currently under study, including continuous quantum walks modeling transport phenomena. Additionally, it proposes zero forcing as a new notion in the analysis of complex networks.

## Index Terms

zero forcing, control, quantum system, walk matrix, Lie algebra, graph

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I. INTRODUCTION

14

15 This paper deals with several concepts from different fields such as linear algebra, graph  
 16 theory and quantum and classical (linear) control theory. In the context of dynamics and control  
 17 of systems on networks, it establishes a connection between a notion in graph theory (zero  
 18 forcing) and concepts in control theory (quantum and classical controllability). We review these  
 19 different concepts before we introduce the technical content of the paper and give physical  
 20 motivation for our study.

21 *A. Background*

22 For a dynamical system with a *control input*, the property of *controllability* describes to what  
 23 extent one can go from one state to another with the evolution corresponding to an appropriate  
 24 choice of the control. If all the possible state transfers can be obtained within a natural set (the  
 25 phase space), then the system is said to be *controllable*.

26 For several classes of systems, controllability has been described in detail and controllability  
 27 tests are known. In particular, for *linear systems*

$$28 \quad \dot{\mathbf{x}} = A\mathbf{x} + \sum_{j=1}^s \mathbf{b}_j u_j, \quad (1)$$

29  $A \in \mathbb{R}^{n \times n}$ ,  $\mathbf{b}_j \in \mathbb{R}^n$ ,  $j = 1, 2, \dots, s$ , where both the state  $\mathbf{x} \in \mathbb{R}^n$  and the control functions  
 30  $u_j = u_j(t)$  enter the right hand side linearly, several equivalent conditions of controllability are  
 31 known. The classical controllability condition (see, e.g., [15]) says the system (1) is controllable  
 32 if and only if the  $n \times (ns)$  matrix

$$33 \quad \widetilde{W}(A, B) := [\mathbf{b}_1, A\mathbf{b}_1, \dots, A^{n-1}\mathbf{b}_1, \dots, \mathbf{b}_s, A\mathbf{b}_s, \dots, A^{n-1}\mathbf{b}_s],$$

34 has full rank  $n$ , where  $B := [\mathbf{b}_1 \quad \mathbf{b}_2 \quad \dots \quad \mathbf{b}_s]$  (note that  $\widetilde{W}(A, B)$  is obtained from the *con-*  
 35 *trollability matrix*  $\mathcal{C}(A, B) = [B, AB, \dots, A^{n-1}B]$  by a permutation of the columns). In this  
 36 case, for any prescribed state transfer  $\mathbf{x}_0 \rightarrow \mathbf{x}_1 (\in \mathbb{R}^n)$  and interval  $[0, T]$ , there exists a control  
 37  $\mathbf{u}(t) = [u_1, \dots, u_s]^T$  such that the corresponding solution  $\mathbf{x}(t)$  of (1) satisfies  $\mathbf{x}(0) = \mathbf{x}_0$  and  
 38  $\mathbf{x}(T) = \mathbf{x}_1$ . For quantum mechanical systems which are closed (i.e., not interacting with the  
 39 environment) and finite dimensional, one considers the *Schrödinger equation*

$$40 \quad i \frac{d}{dt} |\psi\rangle = H(\mathbf{u}) |\psi\rangle, \quad (2)$$

41 where  $|\psi\rangle \in \mathbb{C}^n$  is the quantum state and the Hamiltonian matrix  $H = H(\mathbf{u})$  is Hermitian  
 42 and depends on a control  $\mathbf{u} = \mathbf{u}(t)$  which in some cases can be assumed to be a switch  
 43 between different Hamiltonians. If (2) is a system linear in the state  $|\psi\rangle$ , the solution of (2) is  
 44  $|\psi(t)\rangle = X(t)|\psi(0)\rangle$  where  $X = X(t)$  is the solution of the *Schrödinger matrix equation*

$$45 \quad i\dot{X} = H(\mathbf{u})X \quad (3)$$

46 with initial condition equal to the  $n \times n$  identity matrix  $I_n$ . Since  $H = H(\mathbf{u})$  is Hermitian for  
 47 every value of  $\mathbf{u}$  and therefore  $-iH$  is skew-Hermitian, the solution of (3) is forced to be unitary  
 48 at every time  $t$ . In this context, the system is called completely controllable if for any unitary  
 49 matrix  $X_f$  in  $SU(n)$ <sup>1</sup> there exists a control function  $\mathbf{u} = \mathbf{u}(t)$  and an interval  $[0, T]$  such that  
 50 the corresponding solution  $X = X(t)$  of (3) satisfies  $X(0) = I_n$  and  $X(T) = X_f$ .

51 At the beginning of the development of the theory of quantum control, it was realized (see  
 52 e.g., [11]) that system (3) has a structure familiar in geometric control theory [13] and therefore  
 53 controllability conditions developed there can be directly applied. In particular, the Lie algebra  
 54 rank condition [14] says that a necessary and sufficient condition for complete controllability  
 55 of system (3) is that the Lie algebra generated by the matrices  $\{iH(\mathbf{u})\}$  (as  $\mathbf{u}$  varies in the set  
 56 of admissible values for the control) is  $su(n)$  or  $u(n)$ .<sup>2</sup> This has given rise to a comprehensive  
 57 approach to quantum control based on the application of techniques of Lie algebras and Lie  
 58 group theory [7].

59 In recent years, there has been considerable interest in the study of control systems, both  
 60 classical and quantum, which are naturally modeled on networks. One direction in this research  
 61 is provided by the literature on ‘*consensus*’ problems where interconnected systems such as in  
 62 robotics [4], which interact in various ways, cooperate to reach a certain desired collective  
 63 behavior [19], [20]. Often one tries to relate the controllability of systems on networks to  
 64 topological or graph theoretic properties of the network. For quantum systems, the nodes of  
 65 the network may represent energy levels or particles which are interacting with each other. For  
 66 these systems, the application of the Lie algebra rank condition to determine controllability can

<sup>1</sup>Following standard notation,  $SU(n)$  is the special unitary group, i.e., the matrix group of  $n \times n$  unitary matrices having determinant 1.

<sup>2</sup>Following standard notation,  $u(n)$  is the Lie algebra of  $n \times n$  skew-Hermitian matrices and  $su(n)$  is the Lie algebra of  $n \times n$  skew-Hermitian matrices with zero trace.

67 become cumbersome and subject to errors when the dimension of the system becomes large.  
 68 It is preferable to have criteria based on graph theoretic properties of the network not only  
 69 because they are typically checked more efficiently but also because they give more insight in  
 70 the dynamics of the system. Work in this direction has been done in [2], [5], [21]. In this context,  
 71 a relevant property of a graph  $G$  and a subset  $S$  of its vertices is the capability of this set to  
 72 ‘infect’ all the vertices of the graph, as explained in the next paragraph.

73 Every graph discussed is simple (no loops or multiple edges), undirected, and has a finite  
 74 nonempty vertex set. Consider a graph  $G$  and color each of its vertices black or white. A vertex  
 75  $v$  is said to *infect*, or *force*, a vertex  $w$  if  $v$  is black,  $w$  is white,  $w$  is a neighbor of  $v$ , and  $w$   
 76 is the only white neighbor of  $v$ . In the case where infection of  $w$  has occurred, we change the  
 77 color of  $w$  to black and continue the iterative procedure. The set  $S$  is called a *zero forcing set*  
 78 if this procedure, starting from a graph where only the vertices in  $S$  are black, leads to a graph  
 79 where *all* vertices are black. An example of a zero forcing (infection) process is shown in Figure  
 1, indicated by arrows; the set of black vertices is a zero forcing set.

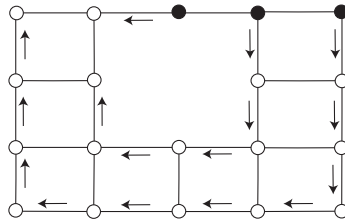


Fig. 1. A zero forcing set and the process by which it can infect all vertices.

80 For a real symmetric  $n \times n$  matrix  $A = [a_{kj}]$ , the *graph* of  $A$ , denoted  $\mathcal{G}(A)$ , is the graph with  
 81 vertices  $\{1, \dots, n\}$  and edges  $\{kj : a_{kj} \neq 0 \text{ and } k \neq j\}$ . Observe that  $G = \mathcal{G}(A_G) = \mathcal{G}(L_G)$ ,  
 82 where  $A_G$  and  $L_G = D_G - A_G$  denote the adjacency matrix of  $G$  and the Laplacian matrix of  
 83  $G$ , respectively (here  $D_G$  is the diagonal matrix of degrees). Zero forcing has been studied in  
 84 detail in the context of linear algebra. This is because the size of a minimum zero forcing set  
 85 of a given graph  $G$ , which is called the *zero forcing number*  $Z(G)$ , is an upper bound to the  
 86 maximum nullity (or maximum co-rank) over any field of  $G$  [3]; the maximum nullity is taken  
 87 over all symmetric matrices  $A$  such that  $\mathcal{G}(A) = G$  (see [8] for background on the problem of  
 88 determining maximum nullity).  
 89

90 Zero forcing appears then to be a valuable concept in the study of graph-theoretic properties  
 91 that are captured by generalized adjacency matrices. Indeed, there are important classical pa-  
 92 rameters introduced with this purpose, e.g., the Colin de Verdière number, the Haemers bound,  
 93 etc. It has to be remarked that questions about the maximum nullity of a graph are generally  
 94 difficult problems and the zero forcing number does not constitute an exception: it was shown  
 95 in [1] that there is no poly-logarithmic approximation algorithm for the zero forcing number.

### 96 *B. Contribution of the paper and physical motivation*

97 In this paper, we consider the dynamics of a system defined on a network and relate the above  
 98 notions and criteria of controllability with the graph theoretic concept of zero forcing. Abstractly,  
 99 we consider a graph  $G$  and a subset  $S = \{j_1, \dots, j_s\}$  of its vertices  $V(G) = \{1, \dots, n\}$ . The  
 100 dynamics are that of a quantum system (3) where the Hamiltonian is allowed to take the values  
 101  $\{A, \mathbf{e}_{j_1} \mathbf{e}_{j_1}^T, \dots, \mathbf{e}_{j_s} \mathbf{e}_{j_s}^T\}$ . Here  $A$  is the adjacency matrix  $A_G$  of  $G$ , Laplacian matrix  $L_G$  of  $G$ ,  
 102 or more generally a real symmetric matrix such that  $\mathcal{G}(A) = G$  with all nonzero off-diagonal  
 103 entries of  $A$  having the same sign (which is the typical situation in transport models). The  
 104 vectors  $\{\mathbf{e}_{j_1}, \dots, \mathbf{e}_{j_s}\}$  are the characteristic vectors<sup>3</sup> of the vertices in  $S$ . In this way, we can  
 105 associate a linear system (1) with  $A$  and  $\mathbf{b}_1 = \mathbf{e}_{j_1}, \dots, \mathbf{b}_s = \mathbf{e}_{j_s}$ . The main result of the present  
 106 paper says that controllability in the quantum sense, expressed by the Lie algebra rank condition,  
 107 and controllability in the sense of linear systems, expressed by the controllability matrix rank  
 108 condition, are equivalent conditions (see Corollary 3.7). Moreover, if the set  $S$  (corresponding  
 109 to  $\mathbf{e}_{j_1}, \dots, \mathbf{e}_{j_s}$ ) is a zero forcing set, then these equivalent controllability conditions are true (see  
 110 Corollary 4.2); the converse is false. The first of these results is along the same lines as the main  
 111 result of [10] which considers the case of quantum dynamics switching between the Hamiltonian  
 112  $A$  and  $\mathbf{z}\mathbf{z}^T$ , where  $\mathbf{z} = \sum_{j \in S} \mathbf{e}_j$ , and establishes the connection between controllability (quantum  
 113 and linear). As mentioned above, these characterizations avoid lengthy calculations of the Lie  
 114 algebra generated by a given set of Hamiltonians and replace them with more easily verified  
 115 graph theoretic and linear algebra tests.

116 On physical grounds, our motivation for considering a Hamiltonian specified by a matrix with  
 117 the given graph comes from the study of *continuous time quantum walks* which model transport

<sup>3</sup>The vector  $\mathbf{e}_j$  has the  $j$ th entry equal to one and every other entry equal to zero and is also called the  $j$ th *standard basis vector*.

118 phenomena in many physical and biological systems [6]. A recent review is given in [18]. Most  
 119 of the studies consider this sole Hamiltonian and concern the statistical (diffusion) properties of  
 120 the dynamics. We add here the Hamiltonians  $e_j e_j^T$  where  $e_j$  is the characteristic vector of a given  
 121 node of the network and study the nature of the states that the resulting dynamics can achieve,  
 122 in particular whether an arbitrary (unitary) state transfer can be achieved between the states of  
 123 the quantum system. The Hamiltonians  $e_j e_j^T$  model a prescribed energy difference between the  
 124 corresponding node and all the other nodes of the network, which are assumed to be at the same  
 125 energy level. Therefore the dynamics is the alternating of a diffusion process (modeled by the  
 126 Hamiltonian  $A$ ) and a rearrangement of the energies of the various states by selecting one of  
 127 the states as high energy state and all the other at the same (lower) energy.

128 Theoretical research in network theory has focused on a number of discrete time, deterministic  
 129 diffusion processes on graphs. While zero forcing has not been studied in this context, there are  
 130 two directions of research that are closely related: as it was already noted in [1], the threshold  
 131 model introduced for studying influence in social networks shares with zero forcing certain issues  
 132 underlying its computational complexity [16]; the model of complex networks controllability  
 133 recently proposed in [17] also makes a natural use of the controllability matrix rank condition  
 134 and it singles out certain combinatorial properties to determine when the condition is satisfied.  
 135 Determining whether zero forcing has a place in the metrology of complex networks is a point  
 136 of further interest.

137 The paper is organized as follows. In Section II we introduce notation and give background and  
 138 basic results concerning Lie algebras that will be used in the following sections. The connection  
 139 between quantum (Lie algebraic) controllability and the controllability matrix rank criterion for  
 140 linear systems is established in Section III. There we also prove the converse of the main result  
 141 of [10]. The relation with the zero forcing property is established in Section IV, and Section V  
 142 contains concluding remarks.

## 143 II. LIE ALGEBRA TERMINOLOGY AND PRELIMINARY RESULTS

144 Standard material on Lie algebras can be found in [12]. For  $A_1, \dots, A_k \in \mathbb{C}^{n \times n}$ ,  $\langle A_1, \dots, A_k \rangle_{[\cdot, \cdot]}$   
 145 denotes the real Lie algebra generated by  $A_1, \dots, A_k$  under addition, real scalar multiplication,  
 146 and the commutator operation. Let  $\mathcal{H}_n(\mathbb{R})$  denote the real vector space of symmetric matrices.  
 147 For  $A \in \mathcal{H}_n(\mathbb{R})$ , the notation  $A = [a_{kj}]$  means for  $k < j$  the  $(k, j)$  and  $(j, k)$  entries of  $A$  are

148 both  $a_{kj}$ . Observe that  $A = [a_{kj}] \in \mathcal{H}_n(\mathbb{R})$  can be expressed as

$$149 \quad A = \sum_{k=1}^n a_{kk} \mathbf{e}_k \mathbf{e}_k^T + \sum_{k < j} a_{kj} (\mathbf{e}_k \mathbf{e}_j^T + \mathbf{e}_j \mathbf{e}_k^T).$$

150 The following proposition is well known (a proof appears in [10]). It provides a link between an  
 151 appropriate Lie algebra of real matrices and the Lie algebra rank condition of quantum control  
 152 theory, thereby allowing us to work with real matrices only. Recall that the Lie algebra consisting  
 153 of all real  $n \times n$  matrices is denoted by  $gl(n, \mathbb{R})$ ,  $sl(n, \mathbb{R})$  denotes the Lie algebra of real  $n \times n$   
 154 matrices with zero trace,  $u(n)$  denotes the Lie algebra of all skew-Hermitian (complex)  $n \times n$   
 155 matrices, and  $su(n)$  denotes the Lie algebra of all skew-Hermitian (complex)  $n \times n$  matrices  
 156 with zero trace. All these Lie algebras are considered as vector spaces over the field of real  
 157 numbers.

158 *Proposition 2.1:* For  $A_1, \dots, A_k \in \mathcal{H}_n(\mathbb{R})$ ,

$$159 \quad \langle A_1, \dots, A_k \rangle_{[\cdot, \cdot]} = gl(n, \mathbb{R}) \iff \langle iA_1, \dots, iA_k \rangle_{[\cdot, \cdot]} = u(n).$$

160 The next lemma is used in the proof of Theorem 3.6 in the next section.

161 *Lemma 2.2:* Let  $A, B_1, \dots, B_s \in \mathcal{H}_n(\mathbb{R})$ , with  $s \geq 1$ . Define  $\mathcal{L} := \langle A, B_1, \dots, B_s \rangle_{[\cdot, \cdot]}$  and let  
 162  $\hat{\mathcal{L}}$  denote the smallest ideal of  $\mathcal{L}$  that contains  $B_i, i = 1, \dots, s$ . If  $\mathcal{L} = gl(n, \mathbb{R})$  and  $\text{tr } B_k \neq 0$   
 163 for some  $B_k$ , then  $\hat{\mathcal{L}} = gl(n, \mathbb{R})$ .

164 *Proof:* For  $n = 1$  the result is clear, so assume  $n \geq 2$ ,  $\mathcal{L} = gl(n, \mathbb{R})$ , and  $\text{tr } B_k \neq 0$  for  
 165 some  $B_k$ . Observe that  $\mathcal{L} := \langle A, B_1, \dots, B_s \rangle_{[\cdot, \cdot]}$  is spanned by  $A$  and  $\hat{\mathcal{L}}$ . Since  $\mathcal{L} = gl(n, \mathbb{R})$ , we  
 166 have

$$167 \quad [gl(n, \mathbb{R}), gl(n, \mathbb{R})] = [\text{span}(A) + \hat{\mathcal{L}}, \text{span}(A) + \hat{\mathcal{L}}] \subseteq \hat{\mathcal{L}}.$$

168 It is known that  $[gl(n, \mathbb{R}), gl(n, \mathbb{R})] = sl(n, \mathbb{R})$ , because  $[gl(n, \mathbb{R}), gl(n, \mathbb{R})]$  is a nonzero ideal in  
 169  $sl(n, \mathbb{R})$  and  $sl(n, \mathbb{R})$  is a simple Lie algebra. Since  $\dim sl(n, \mathbb{R}) = n^2 - 1$  and  $B_k \notin sl(n, \mathbb{R})$ ,  
 170  $\dim \hat{\mathcal{L}} \geq n^2$ . Thus  $\hat{\mathcal{L}} = gl(n, \mathbb{R})$ . ■

171 The next lemma is used in the proof of Theorem 3.1 in the next section. Let  $\mathcal{L}$  be a Lie  
 172 algebra,  $A \in \mathcal{L}$ , and let  $\mathcal{K}$  be a subspace of  $\mathcal{L}$ . Recall that the operation  $ad_A$  is defined as  
 173  $ad_A(B) := [A, B]$ , and the *normalizer* of  $\mathcal{K}$  is

$$174 \quad N_{\mathcal{L}}(\mathcal{K}) = \{A : [A, B] \in \mathcal{K} \text{ for all } B \in \mathcal{K}\}.$$

175 It follows from the Jacobi identity that  $N_{\mathcal{L}}(\mathcal{K})$  is a subalgebra of  $\mathcal{K}$  [12, p. 7].

176 *Lemma 2.3:* Let  $A, L \in \mathcal{H}_n(\mathbb{R})$ . Assume  $\langle iA, iL \rangle_{[\cdot, \cdot]} = u(n)$  and define

$$177 \quad \mathcal{S} := \text{span}(\{ad_{iA}^{k_1} ad_{iL}^{k_2} \cdots ad_{iA}^{k_{s-1}} ad_{iL}^{k_s} [iA, iL]\}), \quad (4)$$

178 where  $s$  and  $k_1, \dots, k_s$  are nonnegative integers. Then  $\mathcal{S} = su(n)$ .

179 *Proof:* First note that  $[iA, iL] \neq 0$  because we have assumed that  $iA$  and  $iL$  generate  $u(n)$ .  
 180 Clearly  $iA, iL \in N_{u(n)}(\mathcal{S})$ . Since  $N_{u(n)}(\mathcal{S})$  is a subalgebra of  $u(n)$  and  $iA$  and  $iL$  generate  
 181  $u(n)$ ,  $N_{u(n)}(\mathcal{S}) = u(n)$ . Thus  $\mathcal{S}$  is an ideal of  $u(n)$ . Notice that  $\mathcal{S} \subseteq su(n)$  since  $[iA, iL]$  is  
 182 skew-Hermitian with zero trace and  $iA$  and  $iL$  are skew-Hermitian. Since  $\mathcal{S}$  is an ideal of  $u(n)$ ,  
 183  $\mathcal{S}$  is an ideal of  $su(n)$ , and  $\mathcal{S} \neq \{0\}$ . Since  $su(n)$  is a simple Lie algebra, by definition it has  
 184 only the trivial ideals  $\{0\}$  and  $su(n)$ . Therefore  $\mathcal{S} = su(n)$ . ■

185 For  $A \in \mathcal{H}_n(\mathbb{R})$  and  $Z = \{\mathbf{z}_1, \dots, \mathbf{z}_s\} \subset \mathbb{R}^n$ , the *real Lie algebra generated by  $A$  and  $Z$*  is  
 186 defined as

$$187 \quad \mathcal{L}(A, Z) := \langle A, \mathbf{z}_1 \mathbf{z}_1^T, \dots, \mathbf{z}_s \mathbf{z}_s^T \rangle_{[\cdot, \cdot]}. \quad (5)$$

### 188 III. CONTROLLABILITY AND WALK MATRICES

189 In this section we show that controllability in the quantum sense, expressed by the Lie algebra  
 190 rank condition, and controllability in the sense of linear systems, expressed by the controllability  
 191 matrix rank condition, are equivalent.

192 For  $A \in \mathcal{H}_n(\mathbb{R})$  and  $Z = \{\mathbf{z}_1, \dots, \mathbf{z}_s\} \subset \mathbb{R}^n$ , the *extended walk matrix* of  $A$  and  $Z$  is the  
 193  $n \times (ns)$  real matrix

$$194 \quad \widetilde{W}(A, Z) := [\mathbf{z}_1, A\mathbf{z}_1, \dots, A^{n-1}\mathbf{z}_1, \dots, \mathbf{z}_s, A\mathbf{z}_s, \dots, A^{n-1}\mathbf{z}_s]. \quad (6)$$

195 A special case is when  $Z = Z_S := \{\mathbf{e}_j : j \in S\}$  (with  $\mathbf{e}_j$  denoting the  $j$ -th standard basis vector)  
 196 for some subset  $S \subseteq V(G)$  for a graph  $G$  and  $A$  is the adjacency matrix  $A_G$  of the graph. In  
 197 this case, the relevant walk matrix is  $\widetilde{W}(A_G, Z_S)$ .

198 For  $s = 1$  the connection between the walk matrix  $\widetilde{W}(A, Z)$  in (6) and the Lie algebra  
 199  $\mathcal{L}(A, Z)$  in (5) was studied in [10]. It was shown [10, Lemma 1] that  $\text{rank } \widetilde{W}(A, \{\mathbf{z}\}) = n$   
 200 implies  $\mathcal{L}(A, \{\mathbf{z}\}) = gl(n, \mathbb{R})$  (although the result is stated for the adjacency matrix and a 0,1-  
 201 vector, the proof remains valid in general), or equivalently,  $\langle iA, i\mathbf{z}\mathbf{z}^T \rangle_{[\cdot, \cdot]} = u(n)$  (cf. Proposition  
 202 2.1). The next theorem states that the converse of this result is also true.



203 *Theorem 3.1:* Consider a matrix  $A$  in  $\mathcal{H}_n(\mathbb{R})$  and a vector  $\mathbf{z} \in \mathbb{R}^{n \times n}$ . Then,  $\langle iA, i\mathbf{z}\mathbf{z}^T \rangle_{[\cdot, \cdot]} =$   
 204  $u(n)$  (or equivalently  $\mathcal{L}(A, \{\mathbf{z}\}) = gl(n, \mathbb{R})$ ) implies that  $\text{rank } \widetilde{W}(A, \{\mathbf{z}\}) = n$ .

205 *Proof:* The equivalence of the hypotheses is justified by Proposition 2.1. The result is clear  
 206 if  $n = 1$ , so assume  $n \geq 2$ . We use a contradiction argument. Assume the rank of the walk  
 207 matrix  $\widetilde{W}(A, \{\mathbf{z}\})$  is less than  $n$  but  $\langle iA, iL \rangle_{[\cdot, \cdot]} = u(n)$ , where  $L := \mathbf{z}\mathbf{z}^T$ . There exists a vector  
 208  $\mathbf{x} \in \mathbb{R}^n$  such that  $\mathbf{x}^* \widetilde{W}(A, \{\mathbf{z}\}) = 0$ . Consider the rank 1 matrix  $D := \mathbf{x}\mathbf{x}^*$ . We claim that  $D$   
 209 commutes with every matrix in  $\mathcal{S}$ , where  $\mathcal{S}$  is as in (4). To see this, notice that from (4), all  
 210 elements in  $\mathcal{S}$  are linear combinations of monomials of the form  $M = A^{k_1} L^{k_2} A^{k_3} \dots L^{k_{p-1}} A^{k_p}$ ,  
 211 for some  $p \geq 1$ ,  $k_j \geq 0$ , and  $L$  appearing at least once with exponent greater than zero. When  
 212 multiplying  $D$  with  $M$ , with  $D$  on the left, write  $M$  as  $A^{k_1} LY$  for some matrix  $Y$ , so we have

$$213 \quad DM = DA^{k_1} LY = \mathbf{x}\mathbf{x}^* A^{k_1} \mathbf{z}\mathbf{z}^* Y = 0, \quad (7)$$

214 which follows immediately from the condition  $\mathbf{x}^* \widetilde{W}(A, \{\mathbf{z}\}) = 0$  for  $n - 1 \geq k_1 \geq 0$ , and by  
 215 using the Cayley-Hamilton theorem for  $k_1 \geq n$ . Analogously, when multiplying  $D$  on the right  
 216 of  $M$ , we write  $M$  as  $QLA^{k_p}$ , for some matrix  $Q$ , and we have

$$217 \quad MD = QLA^{k_p} D = Q\mathbf{z}\mathbf{z}^* A^{k_p} \mathbf{x}\mathbf{x}^* = 0, \quad (8)$$

218 since  $\mathbf{x}^* A^{k_p} \mathbf{z} = 0$  also implies  $\mathbf{z}^* A^{k_p} \mathbf{x} = 0$ . Therefore  $D$  commutes with all elements of  $\mathcal{S}$ .

219 Observe that since  $su(n)$  is simple,  $su(n)$  is an irreducible representation of  $su(n)$ . Therefore,  
 220 since  $D$  commutes with all elements of  $\mathcal{S}$ , it follows from Schur's Lemma that  $D$  must be a  
 221 scalar multiple of the identity [12, p. 26]. However this is not possible since  $D$  has rank 1. This  
 222 gives the desired contradiction and thus completes the proof. ■

223 We study the generalization of this result to multiple vectors ( $s \geq 1$ ) but for matrices  $A$   
 224 and vectors  $\mathbf{z}_1, \dots, \mathbf{z}_s$  related to a connected graph  $G$ . In particular,  $\mathcal{G}(A) = G$ , all nonzero  
 225 off-diagonal entries of  $A$  have the same sign, and  $\mathbf{e}_{j_1}, \dots, \mathbf{e}_{j_s}$  will be the characteristic vectors  
 226 associated to a subset  $S$  of the vertices. In the next section we will relate this to the zero forcing  
 227 property of the set  $S$ . In the context of graphs, it is important to consider multiple vectors  
 228 because if  $G$  is a graph and  $\text{rank } A_G \leq |G| - 2$ , then  $\text{rank } \widetilde{W}(A_G, \{\mathbf{z}\}) < n$  for any one vector  
 229  $\mathbf{z}$ . On the other hand we will see that if  $S$  is a zero forcing set for  $G$  and  $\mathcal{G}(A) = G$ , then  
 230  $\mathcal{L}(A, \{\mathbf{e}_j : j \in S\}) = \mathcal{H}_n(\mathbb{R})$  (see Theorem 4.1 below).

231 The next definition extends the definition given in [9] (and implicitly in [10]) of an associative  
 232 algebra that links the walk matrix and controllability. For  $A \in \mathcal{H}_n(\mathbb{R})$  and  $Z = \{\mathbf{z}_1, \dots, \mathbf{z}_s\} \subset$

233  $\mathbb{R}^n$ , define

$$234 \quad P(A, Z) := \{A^m \mathbf{z}_k \mathbf{z}_j^T A^\ell : 1 \leq k, j \leq s, 0 \leq m, \ell \leq n-1\}.$$

235 *Remark 3.2:* For  $A \in \mathcal{H}_n(\mathbb{R})$  and  $Z = \{\mathbf{z}_1, \dots, \mathbf{z}_s\} \subset \mathbb{R}^n$ , the associative algebra generated  
236 by  $P(A, Z)$  is equal to  $\text{span } P(A, Z)$ , because

$$237 \quad (A^m \mathbf{z}_k \mathbf{z}_j^T A^\ell)(A^g \mathbf{z}_p \mathbf{z}_q^T A^h) = (\mathbf{z}_j^T A^{\ell+g} \mathbf{z}_p) A^m \mathbf{z}_k \mathbf{z}_q^T A^h \text{ and } \mathbf{z}_j^T A^{\ell+g} \mathbf{z}_p \in \mathbb{R}.$$

238 *Lemma 3.3:* For  $A \in \mathcal{H}_n(\mathbb{R})$  and  $Z = \{\mathbf{z}_1, \dots, \mathbf{z}_s\} \subset \mathbb{R}^n$ ,  $\text{rank } \widetilde{W}(A, Z) = n$  if and only if  
239  $\text{span } P(A, Z) = \mathbb{R}^{n \times n}$ .<sup>4</sup>

240 *Proof:* Clearly  $\text{rank } \widetilde{W}(A, Z) = n$  if and only if  $\text{range } \widetilde{W}(A, Z) = \mathbb{R}^n$ . First assume  
241  $\text{rank } \widetilde{W}(A, Z) = n$ . For any matrix  $M \in \mathbb{R}^{n \times n}$  with  $\text{rank } M = r$ , there exist vectors  $\mathbf{x}^{(q)}, \mathbf{y}^{(q)}$ ,  
242  $q = 1, \dots, r$ , such that  $M = \sum_{q=1}^r \mathbf{x}^{(q)} \mathbf{y}^{(q)T}$ . Since  $\text{range } \widetilde{W}(A, Z) = \mathbb{R}^n$ , each  $\mathbf{x}^{(q)}$  is expressible  
243 as a linear combination of the columns of  $\widetilde{W}(A, Z)$ , i.e., as a linear combination of vectors of  
244 the form  $A^m \mathbf{z}_k$ , and similarly for  $\mathbf{y}^{(q)}$ . Thus each  $\mathbf{x}^{(q)} \mathbf{y}^{(q)T}$ , and hence  $M$ , is expressible as a  
245 linear combination of  $A^m \mathbf{z}_k \mathbf{z}_j^T A^\ell$ . Thus the matrices of the form  $A^m \mathbf{z}_k \mathbf{z}_j^T A^\ell$  span  $\mathbb{R}^{n \times n}$ .

246 For the converse, observe that if  $B = \{\mathbf{b}_1, \dots, \mathbf{b}_r\}$  is a basis for  $\text{range } \widetilde{W}(A, Z)$ , then

$$247 \quad \text{span } P(A, Z) = \text{span}(\{\mathbf{b}_k \mathbf{b}_j^T : 1 \leq k, j \leq r\}).$$

248 If  $n > r = \text{rank } \widetilde{W}(A, Z)$ , then  $\dim \text{span } P(A, Z) \leq r^2 < n^2 = \dim \mathbb{R}^{n \times n}$ , so the matrices in  
249  $P(A, Z)$  cannot span  $\mathbb{R}^{n \times n}$ . ■

250 The *distance* between two distinct vertices  $u$  and  $v$  of a connected graph  $G$ , denoted by  $d(u, v)$   
251 is the minimum number of edges in a path from  $u$  to  $v$ .

252 *Lemma 3.4:* Let  $A \in \mathcal{H}_n(\mathbb{R})$  such that  $\mathcal{G}(A)$  is connected and all nonzero off-diagonal entries  
253 of  $A$  have the same sign. If  $k, j \in \{1, \dots, n\}$  and  $k \neq j$ , then  $(A^{d(k,j)})_{kj} \neq 0$ .

254 *Proof:* Let  $d := d(k, j)$ . The entry  $(A^d)_{kj}$  is a sum of terms each of which is the product  
255 of  $d$  nonzero entries of  $A$ . Since  $d$  is the distance between  $k$  and  $j$ , only off-diagonal entries  
256 can appear in this product. Thus every term has the same sign and  $(A^d)_{kj} \neq 0$ . ■

257 *Lemma 3.5:* Let  $A \in \mathcal{H}_n(\mathbb{R})$  be such that  $\mathcal{G}(A)$  is connected and all nonzero off-diagonal  
258 entries of  $A$  have the same sign. Let  $S \subseteq \{1, \dots, n\}$  and  $Z = \{\mathbf{e}_j : j \in S\}$  be the subset of  
259 standard basis vectors. Then  $\text{span } P(A, Z) \subseteq \mathcal{L}(A, Z)$ .

<sup>4</sup>As a vector space,  $\mathbb{R}^{n \times n}$  is the same as  $gl(n, \mathbb{R})$ . We use the latter notation when we want to stress the Lie algebra structure on  $gl(n, \mathbb{R})$ .

260 *Proof:* The proof of Lemma 1 in [10] shows that for any real symmetric matrix  $A$  and  
 261 vector  $\mathbf{z}$ ,  $A^m \mathbf{z} \mathbf{z}^T A^\ell \in \mathcal{L}(A, \{\mathbf{z}\})$  for all  $m, \ell \in \{0, \dots, n-1\}$ . Applying this, we obtain that  
 262  $A^m \mathbf{e}_j \mathbf{e}_j^T A^\ell \in \mathcal{L}(A, Z)$  for all  $j \in \{1, \dots, s\}, m, \ell \in \{0, \dots, n-1\}$ . The result will follow if  
 263 we are able to show that  $A^m \mathbf{e}_k \mathbf{e}_j^T A^\ell \in \mathcal{L}(A, Z)$  for all  $k, j \in \{1, \dots, s\}, m, \ell \in \{0, \dots, n-1\}$ ,  
 264 with  $k$  different from  $j$ .

265 Consider the distance  $d := d(k, j)$  between the nodes  $k$  and  $j$  in  $\mathcal{G}(A)$ , which is  $\leq n-1$   
 266 because  $\mathcal{G}(A)$  is connected. From the fact that both  $\mathbf{e}_k \mathbf{e}_k^T$  and  $A^d \mathbf{e}_j \mathbf{e}_j^T$  are in  $\mathcal{L}(A, Z)$ , we have  
 267 in  $\mathcal{L}(A, Z)$ ,

$$268 \quad [\mathbf{e}_k \mathbf{e}_k^T, A^d \mathbf{e}_j \mathbf{e}_j^T] = \mathbf{e}_k \mathbf{e}_k^T A^d \mathbf{e}_j \mathbf{e}_j^T - A^d \mathbf{e}_j \mathbf{e}_j^T \mathbf{e}_k \mathbf{e}_k^T = (\mathbf{e}_k^T A^d \mathbf{e}_j) \mathbf{e}_k \mathbf{e}_j^T.$$

269 It follows from Lemma 3.4 that  $\mathbf{e}_k^T A^d \mathbf{e}_j = (A^d)_{kj} \neq 0$ , and so  $\mathbf{e}_k \mathbf{e}_j^T \in \mathcal{L}(A, Z)$ .

270 Then

$$271 \quad [A^m \mathbf{e}_k \mathbf{e}_k^T, \mathbf{e}_k \mathbf{e}_j^T] = A^m \mathbf{e}_k \mathbf{e}_k^T \mathbf{e}_k \mathbf{e}_j^T - \mathbf{e}_k \mathbf{e}_j^T A^m \mathbf{e}_k \mathbf{e}_k^T \\ 272 \quad = A^m \mathbf{e}_k \mathbf{e}_j^T - (\mathbf{e}_j^T A^m \mathbf{e}_k) \mathbf{e}_k \mathbf{e}_k^T.$$

273 So,  $A^m \mathbf{e}_k \mathbf{e}_j^T \in \mathcal{L}(A, Z)$ . Similarly,  $\mathbf{e}_k \mathbf{e}_j^T A^\ell \in \mathcal{L}(A, Z)$ . Finally,

$$274 \quad [A^m \mathbf{e}_k \mathbf{e}_k^T, \mathbf{e}_k \mathbf{e}_j^T A^\ell] = A^m \mathbf{e}_k \mathbf{e}_k^T \mathbf{e}_k \mathbf{e}_j^T A^\ell - \mathbf{e}_k \mathbf{e}_j^T A^{m+\ell} \mathbf{e}_k \mathbf{e}_k^T \\ 275 \quad = A^m \mathbf{e}_k \mathbf{e}_j^T A^\ell - (\mathbf{e}_j^T A^{m+\ell} \mathbf{e}_k) \mathbf{e}_k \mathbf{e}_k^T.$$

276 So,  $A^m \mathbf{e}_k \mathbf{e}_j^T A^\ell \in \mathcal{L}(A, Z)$ . ■

277 The following theorem establishes the connection between quantum Lie algebraic controlla-  
 278 bility and the rank condition for an extended walk matrix modeled on a graph.

279 *Theorem 3.6:* Let  $A \in \mathcal{H}_n(\mathbb{R})$  such that  $\mathcal{G}(A)$  is connected and all the nonzero off-diagonal  
 280 elements of  $A$  have the same sign. Let  $S \subseteq \{1, \dots, n\}$  and  $Z = \{\mathbf{e}_j : j \in S\}$  be a subset of  
 281 standard basis vectors. Then  $\text{rank } \widetilde{W}(A, Z) = n$  if and only if  $\mathcal{L}(A, Z) = gl(n, \mathbb{R})$ .

282 *Proof:* By Lemma 3.3,  $\text{span } P(A, Z) = \mathbb{R}^{n \times n}$  if and only if  $\text{rank } \widetilde{W}(A, Z) = n$ , so it  
 283 suffices to show that  $\text{span } P(A, Z) = \mathbb{R}^{n \times n}$  if and only if  $\mathcal{L}(A, Z) = gl(n, \mathbb{R})$ . By Lemma  
 284 3.5,  $\text{span } P(A, Z) \subseteq \mathcal{L}(A, Z)$ , so  $\text{span } P(A, Z) = \mathbb{R}^{n \times n}$  implies  $\mathcal{L}(A, Z) = gl(n, \mathbb{R})$ . For the  
 285 converse, assume  $\mathcal{L}(A, Z) = gl(n, \mathbb{R})$ . Then, by Lemma 2.2,  $\hat{\mathcal{L}} = gl(n, \mathbb{R})$ , where  $\hat{\mathcal{L}}$  is the  
 286 smallest ideal of  $\mathcal{L}(A, Z)$  that contains  $\mathbf{e}_j \mathbf{e}_j^T, j = 1, \dots, s$ . It is clear that  $\hat{\mathcal{L}} \subseteq \text{span } P(A, Z)$ ,  
 287 so  $\text{span } P(A, Z) = \mathbb{R}^{n \times n}$ . ■

288 *Corollary 3.7:* Let  $A \in \mathcal{H}_n(\mathbb{R})$  such that  $\mathcal{G}(A)$  is connected and all the nonzero off-diagonal  
 289 elements of  $A$  have the same sign, and let  $S \subseteq \{1, \dots, n\}$ . Then  $\text{rank } \widetilde{W}(A, \{\mathbf{e}_j : j \in S\}) = n$   
 290 if and only if  $\langle iA, \{i\mathbf{e}_j\mathbf{e}_j^T : j \in S\} \rangle_{[\cdot, \cdot]} = u(n)$ , i.e., the quantum system associated with the  
 291 Hamiltonians  $iA$  and  $i\mathbf{e}_j\mathbf{e}_j^T$ ,  $j = 1, \dots, s$ , is controllable.

292 Observe that for any connected graph  $G$ , the adjacency matrix  $A_G$  and the Laplacian matrix  
 293  $L_G$  satisfy the hypotheses of Theorem 3.6 and Corollary 3.7.

294 The result of [10] for the case  $s = 1$  showing that  $\text{rank } \widetilde{W}(A, \{\mathbf{z}\}) = n$  implies  $\langle iA, i\mathbf{z}\mathbf{z}^T \rangle_{[\cdot, \cdot]} =$   
 295  $u(n)$ , (and the converse proved in Theorem 3.1 in this paper) were proved in reference to systems  
 296 on graphs. The proofs however go through for an arbitrary symmetric matrix  $A$  and vector  $\mathbf{z}$ . It  
 297 is natural to ask whether the conditions on the matrix  $A$  that we have used in Theorem 3.6 are  
 298 really necessary. To this purpose, we can observe that the result is not true if we give up either  
 299 of the hypotheses that 1)  $\mathcal{G}(A)$  is connected or 2) the off-diagonal entries of  $A$  have the same  
 300 sign, as shown in the next two examples.

301 *Example 3.8:* To see the necessity of assuming that  $\mathcal{G}(A)$  is connected, consider a block  
 302 diagonal matrix  $A = \begin{bmatrix} A_1 & 0 \\ 0 & A_2 \end{bmatrix}$  with  $A_1$  and  $A_2$  symmetric matrices of dimensions  $n_1$  and  $n_2$ ,  
 303 respectively, with  $n_1 + n_2 = n$ , and  $\mathbf{z}_1$  and  $\mathbf{z}_2$  two vectors that have zeros in the last  $n_2$  or first  
 304  $n_1$  entries, respectively, and such that the corresponding matrices  $\widetilde{W}(A_1, \{\mathbf{z}_1\})$  and  $\widetilde{W}(A_2, \{\mathbf{z}_2\})$   
 305 have ranks  $n_1$  and  $n_2$ , respectively. In this case, the walk matrix  $\widetilde{W}(A, \{\mathbf{z}_1, \mathbf{z}_2\})$  has rank  $n$ , but  
 306 the Lie algebra generated by  $A$ ,  $\mathbf{z}_1\mathbf{z}_1^T$ , and  $\mathbf{z}_2\mathbf{z}_2^T$  contains only block diagonal matrices.

307 *Example 3.9:* To see the necessity of assuming that all nonzero off-diagonal entries of  $A$  have  
 308 the same sign, consider  $A = \begin{bmatrix} 0 & 1 & 0 & 1 \\ 1 & 0 & -1 & 0 \\ 0 & -1 & 0 & 1 \\ 1 & 0 & 1 & 0 \end{bmatrix}$ , and  $Z = \{\mathbf{e}_1, \mathbf{e}_3\}$ . It is straightforward to verify  
 309 that the walk matrix  $\widetilde{W}(A, \{\mathbf{e}_1, \mathbf{e}_3\})$  has rank 4. However,  $\dim \mathcal{L}(A, \{\mathbf{e}_1, \mathbf{e}_3\}) \leq 8$ , as can be  
 310 seen as follows. Let

311

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$$\mathcal{L} := \text{span}\left( \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}, \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}, \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}, \begin{bmatrix} 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{bmatrix}, \right.$$

$$\left. \begin{bmatrix} 0 & 1 & 0 & 1 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}, \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & -1 \\ 0 & 0 & 0 & 0 \end{bmatrix}, \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & -1 & 0 \end{bmatrix}, \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{bmatrix} \right).$$

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314

Since  $[B, C] \in \mathcal{L}$  for all  $B, C \in \mathcal{L}$ ,  $\mathcal{L}$  is a Lie subalgebra of  $gl(4, \mathbb{R})$ . Clearly  $\dim \mathcal{L} \leq 8$  and

315

$$\mathcal{L}(A, \{\mathbf{e}_1, \mathbf{e}_3\}) \subseteq \mathcal{L}.$$

316

#### IV. ZERO FORCING AND CONTROLLABILITY

317

In this section we show that if the set  $S = \{j_1, \dots, j_s\}$  (corresponding to  $\mathbf{e}_{j_1}, \dots, \mathbf{e}_{j_s}$ ) is a zero forcing set, then the equivalent controllability conditions are true (the converse is false).

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The *neighborhood* of  $v \in V(G)$  is  $N(v) = \{w \in V(G) : w \text{ is adjacent to } v\}$ .

320

*Theorem 4.1:* Let  $A \in \mathcal{H}_n(\mathbb{R})$  such that  $\mathcal{G}(A)$  is connected and all the nonzero off-diagonal entries of  $A$  have the same sign. Let  $V := \{1, 2, \dots, n\}$  be the set of vertices for  $\mathcal{G}(A)$ , and  $S \subseteq V$  be a zero forcing set of  $\mathcal{G}(A)$ . Then

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323

$$\mathcal{L}(A, \{\mathbf{e}_j \mathbf{e}_j^T : j \in S\}) = gl(n, \mathbb{R}).$$

324

*Proof:* After a (possibly empty) sequence of forces, denote by  $T$  the set of currently black vertices, and assume that for all  $k \in T$ ,  $\mathbf{e}_k \mathbf{e}_k^T \in \mathcal{L} := \mathcal{L}(A, \{\mathbf{e}_j \mathbf{e}_j^T : j \in S\})$ . The hypotheses of Lemma 3.5 are satisfied for  $Z = \{\mathbf{e}_j : j \in T\}$ , so for all  $k, \ell \in T$ ,  $\mathbf{e}_k \mathbf{e}_\ell^T \in \mathcal{L}$ .

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If  $T \neq V$ , then there is a vertex  $u \in T$  that has a unique neighbor  $w$  outside  $T$ . For that  $u$  we have

328

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$$[\mathbf{e}_u \mathbf{e}_u^T, A] = \sum_{m \in N(u)} \tilde{a}_{um} (\mathbf{e}_u \mathbf{e}_m^T - \mathbf{e}_m \mathbf{e}_u^T),$$

330

where  $\tilde{a}_{um} := a_{um}$  if  $u < m$  and  $\tilde{a}_{um} := a_{mu}$  if  $m < u$ . For all  $m \in N(u)$  such that  $m \neq w$ ,  $m \in T$ , so  $\mathbf{e}_u \mathbf{e}_m^T, \mathbf{e}_m \mathbf{e}_u^T \in \mathcal{L}$ . Thus  $\mathbf{e}_u \mathbf{e}_w^T - \mathbf{e}_w \mathbf{e}_u^T \in \mathcal{L}$ . Since

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$$[\mathbf{e}_u \mathbf{e}_u^T, \mathbf{e}_u \mathbf{e}_w^T - \mathbf{e}_w \mathbf{e}_u^T] = \mathbf{e}_u \mathbf{e}_w^T + \mathbf{e}_w \mathbf{e}_u^T,$$

333  $\mathbf{e}_w \mathbf{e}_u^T, \mathbf{e}_w \mathbf{e}_u^T \in \mathcal{L}$ . Then

334 
$$[\mathbf{e}_w \mathbf{e}_u^T, \mathbf{e}_u \mathbf{e}_w^T] = \mathbf{e}_w \mathbf{e}_w^T - \mathbf{e}_u \mathbf{e}_u^T$$

335 so  $\mathbf{e}_w \mathbf{e}_w^T \in \mathcal{L}$ . Since  $S$  is a zero forcing set, we obtain  $\mathbf{e}_\ell \mathbf{e}_m^T \in \mathcal{L}$  for all  $\ell, m \in V = \{1, \dots, n\}$ ,  
 336 and thus we conclude that  $\mathcal{L} = gl(n, \mathbb{R})$ . ■

337 Applying Proposition 2.1 we obtain the next corollary.

338 *Corollary 4.2:* If  $G$  is a connected graph,  $A \in \mathcal{H}_n(\mathbb{R})$ , all the nonzero off-diagonal entries of  $A$   
 339 have the same sign, and  $S \subseteq V$  is a zero forcing set of  $G$ , then  $\langle iA, \{\mathbf{e}_j \mathbf{e}_j^T : j \in S\} \rangle_{[\cdot, \cdot]} = u(n)$   
 340 and the corresponding quantum system is controllable.

341 Note that the converse of Theorem 4.1 is false.

342 *Example 4.3:* Consider the path on four vertices  $P_4$  with the vertices numbered in order. The  
 343 set  $\{\mathbf{e}_2\}$  is not a zero forcing set for  $P_4$ . However,

344 
$$\widetilde{W}(A_{P_4}, \{\mathbf{e}_2\}) = \begin{bmatrix} 0 & 1 & 0 & 2 \\ 1 & 0 & 2 & 0 \\ 0 & 1 & 0 & 3 \\ 0 & 0 & 1 & 0 \end{bmatrix} \quad \text{and} \quad \text{rank } \widetilde{W}(A_{P_4}, \{\mathbf{e}_2\}) = 4,$$

345 so  $\mathcal{L}(A_{P_4}, \{\mathbf{e}_2\}) = gl(n, \mathbb{R})$  by Theorem 3.6.

## 346 V. CONCLUSION

347 Motivated by the control and dynamics of systems modeled on networks both classical and  
 348 quantum, we have established a connection between various tests of controllability and the notion  
 349 of zero forcing in graph theory. Lie algebraic quantum controllability is necessary and sufficient  
 350 for (classical) linear controllability of an associated system and both notions are implied by the  
 351 zero forcing property of the associated set of vertices. Linear systems have a very well developed  
 352 theory [15] and it is an open question to investigate to what extent this analogy can be further  
 353 used to discover properties of quantum systems and systems on networks.

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